Electronic Properties of *h*-WO₃ and CuWO₄ Nanocrystals as Determined from X-ray Spectroscopy and First-Principles Band-Structure Calculations

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Abstract—The electronic structure of hexagonal WO₃ and triclinic CuWO₄ nanocrystals, prospective materials for functional electronic devices, has been studied using the X-ray photoelectron spectroscopy (XPS) and X-ray emission spectroscopy (XES) methods. The present XPS and XES results render that the W 5d- and O 2p-like states contribute throughout the whole valence-band region of the *h*-WO₃ and CuWO₄ nanocrystalline materials, however maximum contributions of the O 2p-like states occur in the upper, whilst the W 5d-like states in the lower portions of the valence band, respectively. Calculations of band dispersions for several symmetry directions of CuWO₄ have been made using the first-principles self-consistent full potential linearized augmented plane wave (FP-LAPW) method. The FP-LAPW calculations render that CuWO₄ is an indirect-gap material.

Index Terms—tungstate; nanocrystal; oxide; electronic structure.

I. INTRODUCTION

Tungsten oxides are recognized among effective materials for photocatalytic water splitting with generation of hydrogen and oxygen [1]-[4]. Numerous efforts were spent to prepare WO₃ and complex tungstates crystals as nanobelts, whiskers, nanorods and nanoslabs with potentials for using in nanoarchitecture [2], [5]-[9]. Tungstate crystals and films possess interesting optical, electrochromic, gas-sensing properties and are widely used as functional elements in various device structures [10]-[13]. Complex tungstate crystals typically display a framework-type crystal structure with distorted cavities or channels stuffered with other cations at low-symmetry oxygen environment that results in noncentrosymmetric crystal structure. As a result, many tungstates show nonlinear optical properties or are suitable hosts for lasant ions [14]-[20]. Photocatalytic, optical and

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Yu. M. Solonin is with Frantsevich Institute for Problems of Materials Science, National Academy of Sciences of Ukraine, Kiev, UA-03142, Ukraine (e-mail: solonin@ipms.kiev.ua). electro-physical properties of tungstates are a complex function of phase composition and the electronic structure of the crystals [3], [4], [21]-[23].

Present study is aimed to consider general trends found for the electronic properties of tungsten oxide and complex tungstates. Besides this, low-temperature preparation of tungsten oxide nanocrystals with developed surface area will be proposed. General factors defined chemical stability of top surface of tungstates will be also evaluated. For such a study, nanocrystals of hexagonal tungsten trioxide (h-WO₃) and triclinic copper tungstate (CuWO₄) have been chosen.



Fig. 1. SEM image of *h*-WO₃ nanocrystals.

II. EXPERIMENTAL

Earlier, low-temperature precipitation reaction with nitric acid from aqueous solution of the ammonium salt was developed for fabrication of well-shaped molybdenum and gernanium oxide nanocrystals [24]-[28]. Layered molybdenum and tungsten oxides are the strucrural analogues and possess strongly anisotropic structure. In present study this reaction was tested for the formation of tungsten oxide nanocrystals. The precipitation was produced under $T = 100^{\circ}$ C and pH = 1. As it shown in Fig.1, the *h*-WO₃ nanoplates with 10 nm thick and 100 nm diameter were producted. Phase composition of the final product was confirmed with XRD analysis and Raman spectroscopy [29]. The DTA curve recorded up to $T = 500^{\circ}$ C in air shows the thermal transformation of hexagonal h-WO₃ to monoclinic β -WO₃ at T_t = 390-400°C. Most intensive Raman lines at 967, 685, 223 and 122 cm⁻¹ are attributed to the bridged O-W-O bonding.

A comparative experimental study of the electronic structure of h-WO₃ and CuWO₄ nanocrystals has been made using the XPS and XES methods. Details of sample preparation and of experimental techniques can be found elsewhere [21], [22], [30]-[34].

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III. COMPUTATIONAL DETAILS

Calculations of band dispersions for several symmetry directions of CuWO₄ have been made using the first-principles self-consistent full potential linearized augmented plane wave (FP-LAPW) method with the WIEN97 code [35]. Like in the calculations of total and partial densities of states of copper tungstate [33], in the present FP-LAPW calculations, lattice parameters and positions of the constituent atoms of CuWO₄ have been chosen in accordance with the crystallography data determined for the compound by Kihlborg and Gebert [36]. Crystal structure of CuWO₄ is shown in Fig. 2.



Fig. 2. Crystal structure of CuWO₄. Unit cell is outline. Oxygen lone atoms are omitted for clarity.

The generalized gradient approximation (GGA) by Perdew et al. [37] have been used for calculations of the exchange-correlation potential. The muffin-tin (MT) sphere radii of the constituent atoms of CuWO₄ in the present calculations were assumed to be 0.873 Å for W and O and 1.005 Å for Cu. The $R_{\min}^{MT}k_{\max}$ parameter, where R_{\min}^{MT} denotes the smallest MT sphere radius and k_{max} determines the value of the largest k vector in the plane wave expansion, were assumed to be 10.0 (the charge density was Fourier expanded up to the value $G_{\text{max}} = 14$). In the potential decomposition, the valence wavefunctions inside the MT spheres were expanded up to $l_{\text{max}} = 3$. The basis function of the atomic orbitals of Cu, W, and O was analogous to that employed in [28]. Integration through the Brillouin zone was carried out using the tetrahedron method by Blöchl et al. [38]. The Brillouin zone sampling has been done using 92 k-points within the irreducible part of the zone. The iteration process was checked taking into account the changes of total energy and the calculations were interrupted when for three following iterations the change of total energy was less than $\sim 1 \times 10^{-5}$ eV. Calculations of band dispersions for several symmetry directions of *h*-WO₃ are in progress by our group.

IV. RESULTS AND DISCUSSION

Results of comparison on a common energy scale of the experimental X-ray emission W L β_5 and O K α bands, representing the energy distributions of the W 5d- and O 2p-like states, respectively, for the *h*-WO₃ and CuWO₄ nanocrystals under study are presented in Figs. 3 and 4. In these figures, for comparison, the XPS valence-band spectrum (excited by Mg K α radiation) of the corresponding compound is also presented. The method of matching the X-ray emission bands on a common energy scale was analogous to that applied successfully when studying the electronic structure of transition metal oxides and tungstates [21], [31].

Data of matching the X-ray emission bands on a common energy scale (Figs. 3 and 4) render that, for the h-WO₃ and

CuWO₄ nanocrystals under consideration, the main contributions of the W 5d- and O 2p-like states are observed at the bottom and near the top of the valence band, respectively, with contributions of the mentioned states throughout other portions of the valence band of the compounds. The above experimental results confirm the data of the first-principles FP-LAPW calculations of copper tungstate [33].



Fig. 3. The X-ray emission W L β_5 and O K α bands matched on a common energy scale with the XPS valence-band spectrum of *h*-WO₃ nanocrystals.

Predominant contributions of the W 5d- and O 2p-like states into the lower and upper portions of the valence band, respectively, established for h-WO₃ and CuWO₄, look to be a common peculiarity of the electronic structure of tungsten trioxide and MWO₄ tungstates. In particular, this peculiarity was either experimentally or theoretically detected for a number of MWO_4 (M = Fe, Co, Pb, Cd, Zn, Sr) compounds [39]-[41]. However, main contributions of the valence states associated with the M atom depend significantly on its position in the Periodic Table. Partly, the theoretical band-structure calculations [33] have revealed that the top of the valence band of CuWO₄ is dominated by contributions of Cu 3d-like states. However, Cu states do not contribute significantly into the conduction band of the compound [33]. Additionally, Zhang et al. [40] have established minor contributions of lead atoms into the total density of states (DOS) within the energy region corresponding to the positions of the valence and conduction bands of PbWO₄. Nevertheless, in CaWO₄, calcium atoms contribute dominantly into the upper portion of the conduction band (at energies higher that 7 eV above the top of the valence band of the tungstate [40]). It is worth mentioning that contributions of cadmium atoms dominate the bottom of the valence band CdWO₄ tungstate, however no significant Cd of contributions throughout the conduction band of the compound were detected in calculations by Abraham et al. [41].



Fig. 4. The X-ray emission W L β_3 and O K α bands matched on a common energy scale with the XPS valence-band spectrum of CuWO₄ nanocrystals.

Band dispersions for CuWO₄ are plotted in Fig. 5 for several symmetry directions of the triclinic Brillouin zone (Fig. 6), which is characteristic of the compound under consideration. It is worth mentioning that Fig. 6 was plotted adopting the program [42], and in this figure the coordinates of the *k*-points, within the limited region of the Brillouin zone studied for the band dispersions shown in Fig. 5, are as follows: Γ (0.0 0.0 0.0), X (0.5 0.0 0.0), Y (0.0 0.5 0.0), and Z (0.0 0.0 0.5).



Fig. 5. Electronic bands along selected symmetry paths within the first Brillouin zone of CuWO₄.

From Fig. 5 it is obvious that CuWO₄ is an indirect-gap material. This result confirms the similar suggestion by Ruiz-Fuertes et al. [43] based on measurements of optical-absorption spectra of CuWO₄ up to 18 GPa. From the band dispersions plotted in Fig. 5, we can make the following statement about the band extrema in CuWO₄: the valence band has its maximum at the Γ point at the center of the Brillouin zone, while the minimum of the conduction band is located at the Y point (0.0 0.5 0.0). According to the present FP-LAPW calculations, the value of the indirect gap in CuWO₄ equals to 1.9988 eV. This value is slightly smaller as compared with the experimental value of the energy gap, E_g, of CuWO₄ (E_g = 2.3 eV) reported by Ruiz-Fuertes et al. [43] and that (E_g = 2.06 eV) derived experimentally for this tungstate by Benko et al. [44].

However, it is necessary to mention that first-principles calculations made within the GGA approximation [37] underestimate somewhat the energy gap of semiconductors and insulators. As a result, we can consider that the value E_g = 1.9988 eV obtained theoretically in the present FP-LAPW calculations for CuWO₄ is in good agreement with experimentally measured for this material values E_g = 2.3 eV [43] and E_g = 2.06 eV [44].



Fig. 6. Diagram of the Brillouin zone for a triclinic structure, which is characteristic of CuWO₄, based on Ref. [42].

Some differences in the electronic structure of MWO_4 (M = Ca, Cd, Pb) compounds are obvious in the LAPW

calculations [40] and [41]: the minimum band gap of CaWO₄ is located at the Γ point at the center of the Brillouin zone, while the minimum band gaps do not occur at Γ for cadmium and lead tungstates. As it has been established the above mentioned authors, CdWO₄ reveals the minimum band gap at the Y point, which is located at the center of the Brillouin-zone boundary plane perpendicular to the *b* crystal axis, while PbWO₄ possesses band extrema at wave vectors away from the zone center (in the Σ and Δ directions) with possibly indirect band gaps.

Fig. 5 shows that the valence-band width of CuWO₄ equals about 7.7 eV. This result seems to be in agreement with those obtained in LAPW calculations of electronic properties of tungstates MWO_4 (M = Ca, Cd, Pb) [40] and [41]. Particularly, the valence-band width was evaluated to be about 7 eV for CdWO₄ with the wolframite-type structure [41]. This value is very close to that obtained in the present FP-LAPW calculations of CuWO₄. For scheelite-type CaWO₄ and PbWO₄ compounds the LAPW calculations by Zhang et al. [40] have revealed somewhat smaller magnitudes of the valence-band widths, mainly 5 and 5.5 eV for calcium and lead tungstate, respectively.

The XPS W 4f core-level spectra of the *h*-WO₃ and CuWO₄ nanocrystyals are presented in Fig. 7. It is apparent that the XPS W 4f core-level spectra in the both nanocrystalline materials under consideration are simple spin-doublets with the XPS W $4f_{7/2}$ binding energies corresponding to those of tungsten in the formal valence +6. It should be noted that values of the XPS O 1s core-level binding energies in the nanocrystals studied were found to be close to those in other representative tungsten-bearing wolframite-type [39] as well as KY(WO₄)₂-type [45] tungstates.



Fig. 7. XPS W 4f core-level spectra of h-WO₃ and CuWO₄ nanocrystals.

V. CONCLUSIONS

In summary, the electronic structure of h-WO₃ and CuWO₄ nanocrystals, prospective materials for renewable energy production and functional devices, was studied using X-ray photoelectron spectroscopy (XPS) and X-ray emission spectroscopy (XES). The XPS valence-band and core-level spectra as well as the XES W L β_5 and O K α bands representing the energy distribution of the W 5d- and O 2p-like states, respectively, were measured. Comparison on a common energy scale of the X-ray emission W L β_5 and O K α bands and the XPS valence-band spectra indicates that, the W 5d- and O 2p-like states contribute throughout the whole valence-band region of h-WO₃ and CuWO₄ nanocrystals, however maximum contributions of the O 2p-like states

occur in the upper, whilst the W 5d-like states in the lower portions of the valence band, respectively. Tungsten atoms in the both nanocrystalline materials are in the formal valence +6.

According to the present FP-LAPW calculations, CuWO₄ is an indirect-gap material: the valence band has its maximum at the Γ point at the center of the Brillouin zone, while the minimum of the conduction band is located at the Y point (0.0 0.5 0.0). The FP-LAPW calculations reveal the indirect gap of 1.9988 eV in CuWO₄. This value is in reasonable agreement with the experimental energy gap values of CuWO₄ reported by Ruiz-Fuertes et al. [43] (E_g = 2.3 eV) and by Benko et al. [44] (E_g = 2.06 eV).

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